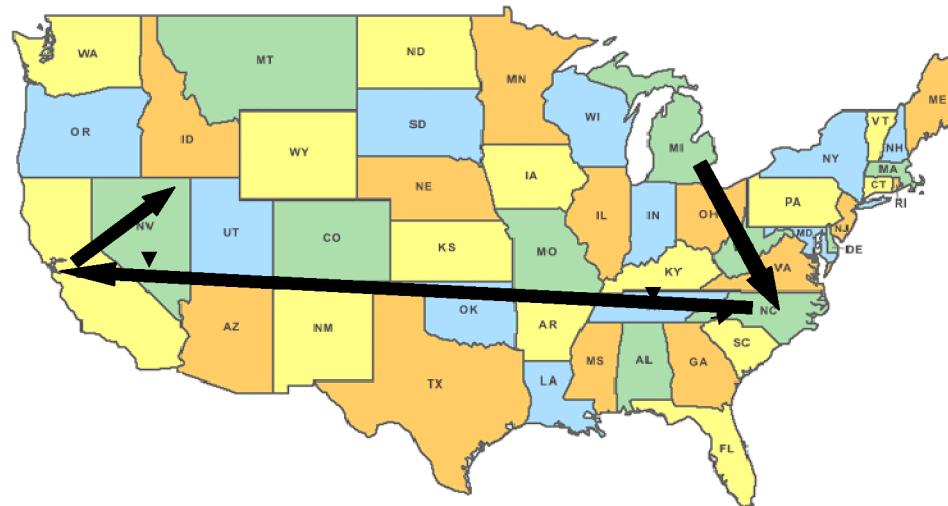


Computational Physics with Particles

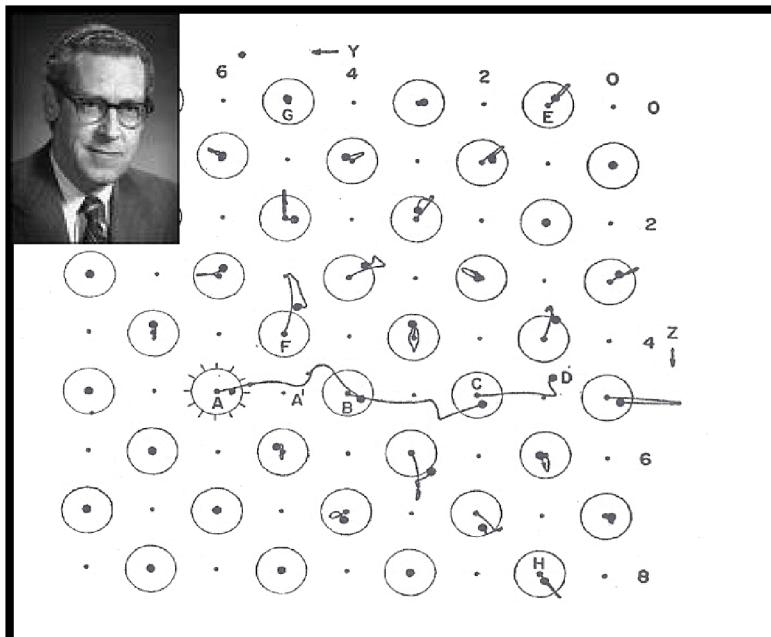
**Wm G Hoover & Carol G Hoover
[no longer at UCDavis & LLNL!]**



**Ruby Valley Research Institute
Highway Contract 60, Box 598
Ruby Valley 89833 Nevada USA**

Historical/AutoBiographical Summary

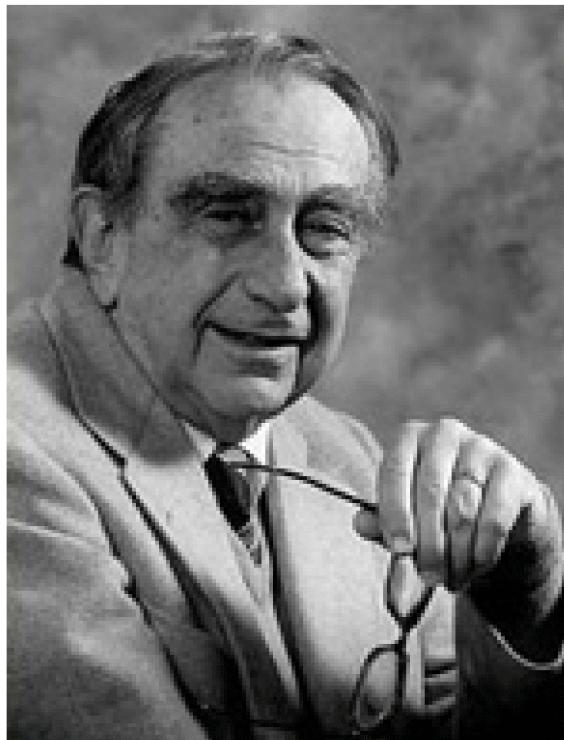
**@ University of Michigan 1958-1961
and saw “Molecules in Motion”
by Berni Alder & Tom Wainwright
Scientific American (1959) .**



Scanned at the American
Institute of Physics

Historical/AutoBiographical Summary

**@ Lawrence Livermore Laboratory &
“Teller Tech” [UCDavis] 1962-2004
[10,000 people/square mile]**



Historical/AutoBiographical Summary

**Bill Ashurst's PhD Thesis 1972-1975 :
“Nonequilibrium Molecular Dynamics”**
**Oyeon Kum's PhD Thesis 1992-1995 :
“Smooth Particle Applied Mechanics”**



Historical/AutoBiographical Summary

Retirement @ Ruby Valley, Nevada 2004

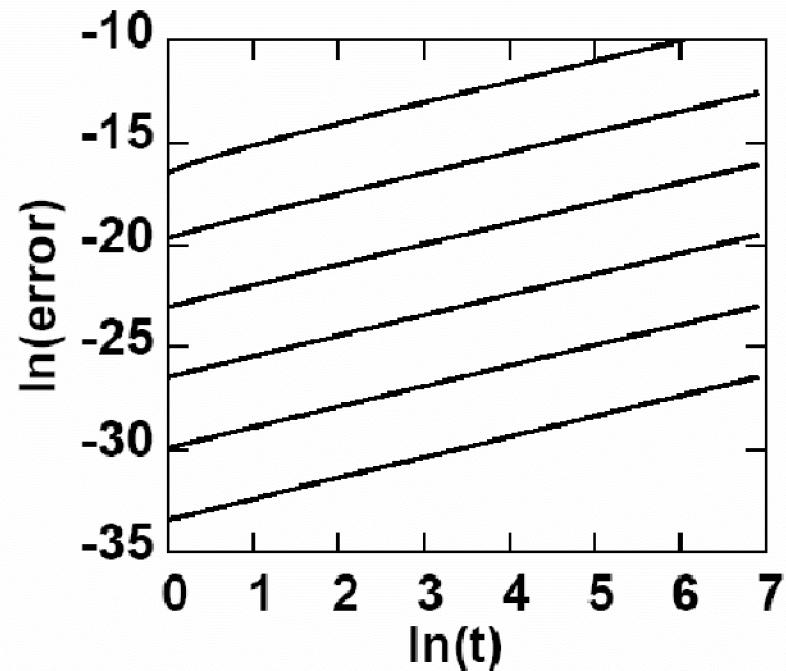


Computational Physics with Particles

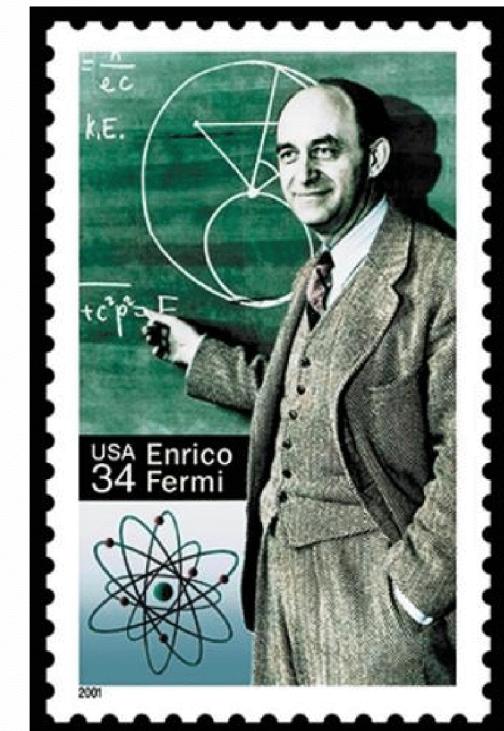
- Particles obey **Ordinary Differential Equations**
- Molecular Dynamics + Control Theory [**NEMD**]
- Continuum Mechanics with Particles [**SPAM**]
- Future Needs for both these Particle Methods

OWOWOWOWOWOWOWOWOWOWOW

Equilibrium Case, using Runge-Kutta



$\Delta t = 0.01 \text{ to } 0.32$



Particles → Ordinary Differential Equations

**Equilibrium & Nonequilibrium Simulations
Need Numerical Solutions *via* Runge-Kutta .**

**Nonequilibrium Simulations require additional
Boundary, Control, and Driving Forces .**

For example, temperature control :

$$m\ddot{r} = F_{\text{Atomistic}} - \zeta p ; \dot{\zeta} \equiv [(K/K_{\text{AVG}}) - 1]/\tau^2 .$$

This generates Gibbs' canonical ensemble .

Nonequilibrium Ordinary Differential Equations

Heat Flow Solutions *via* Runge-Kutta
For **Harmonic** Chain with **Quartic** Tethers

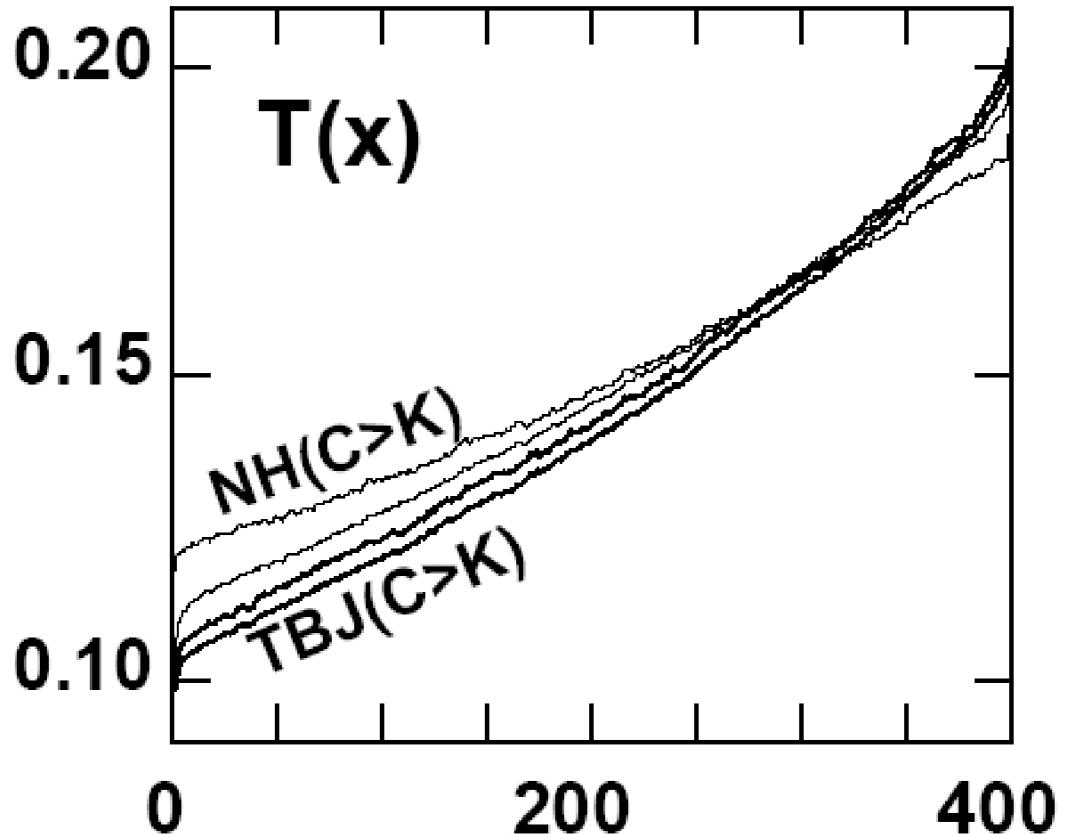
$$\Phi_{\text{Harmonic}} = \frac{1}{2} \sum (|x_{ij}| - d)^2$$

$$\Phi_{\text{Tether}} = \frac{1}{4} \sum (\delta x)^4$$

First (**cold**) and Last (**hot**) Particles
Thermostated with Kinetic or
Configurational Thermostat Forces .

OWOWOWOWOWOWOWOWOWOW

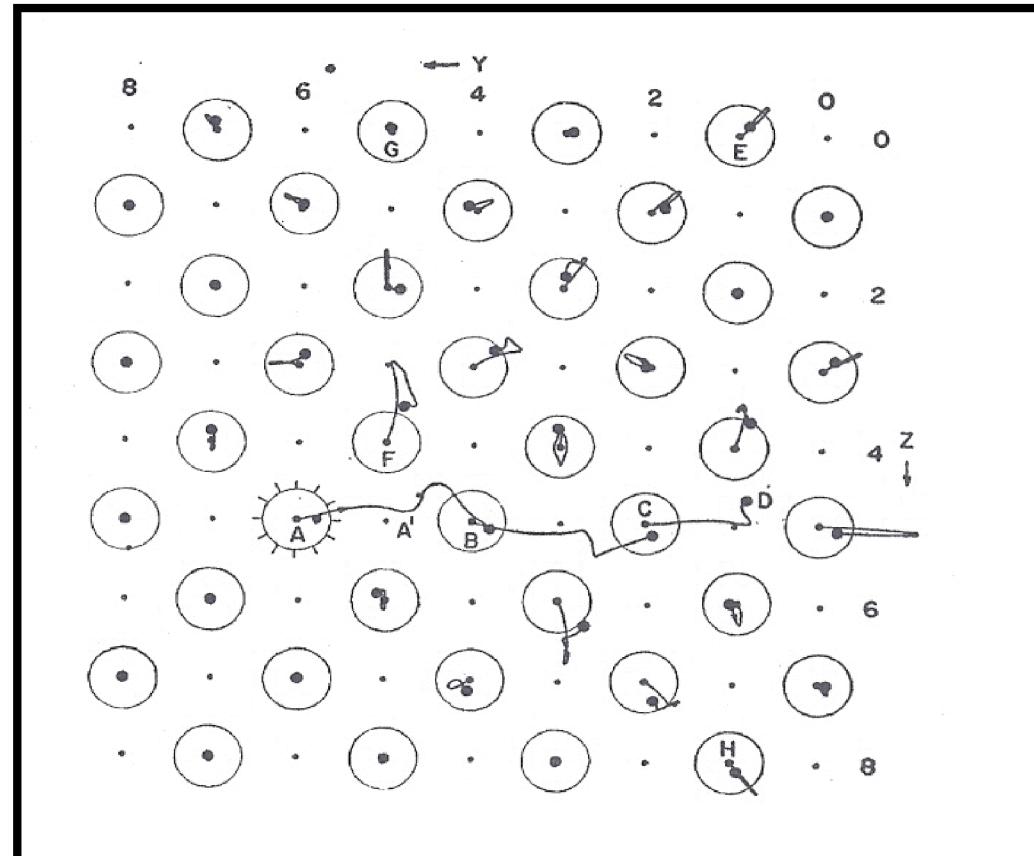
Nonequilibrium Case: Aoki-Kusnezov ϕ^4 model



$$kT_K = \langle p^2/m \rangle \text{ or}$$

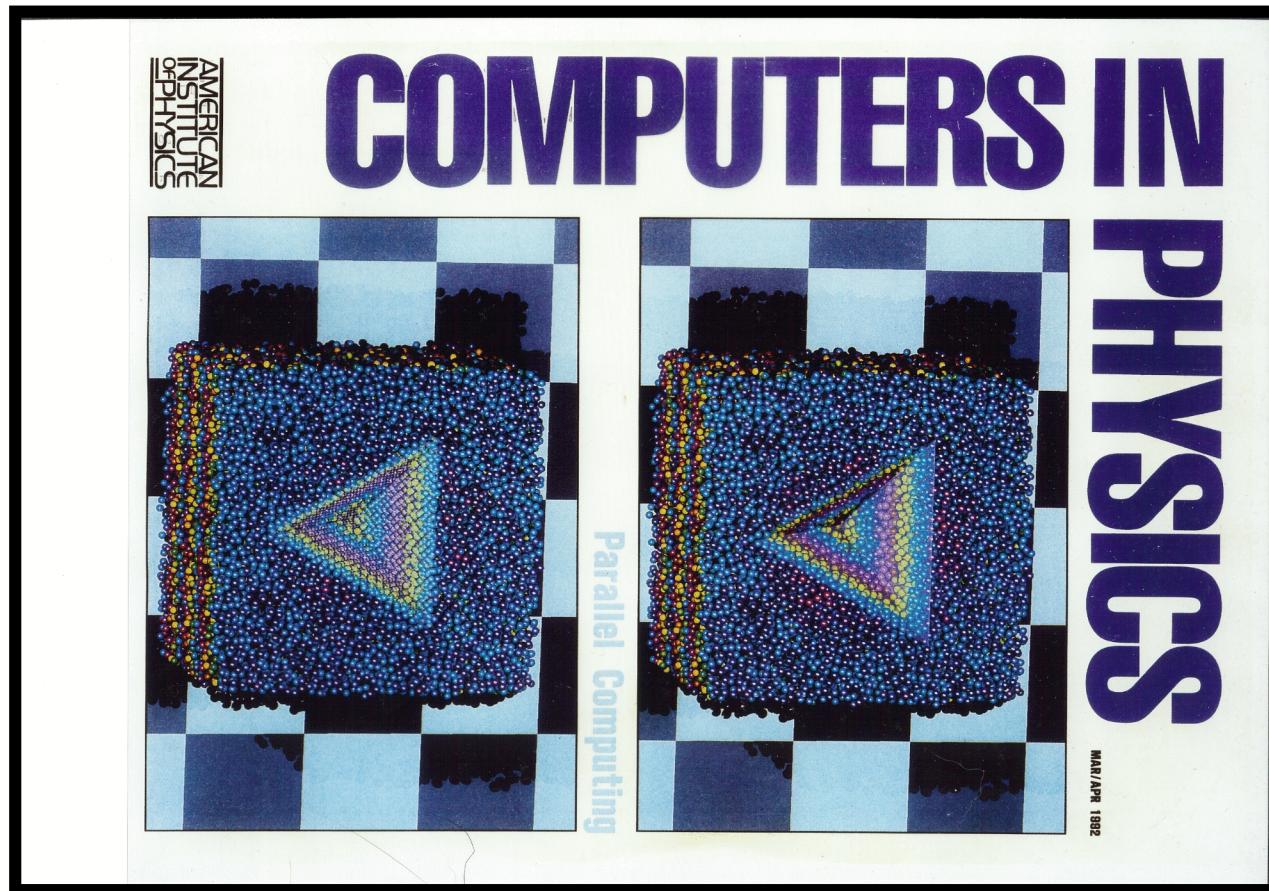
$$kT_C = \langle F^2 \rangle / \langle \nabla^2 \mathcal{H} \rangle$$

1959: Dynamics with 500 Particles



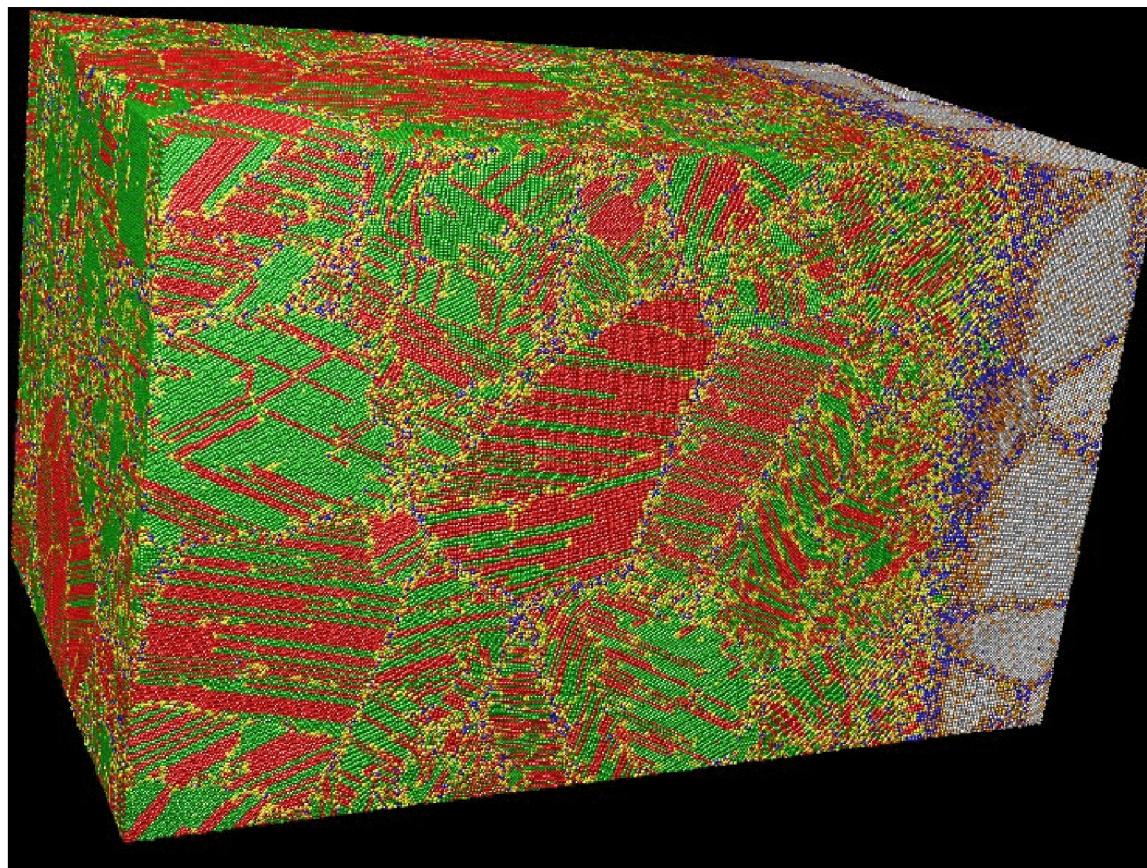
From Vineyard's 1960 Physical Review Article

1989: Dynamics with 1,000,000 Particles



From our work in Japan, 1989-1990, simulating silicon .

2007: Dynamics with 30,000,000 Particles



From Kai Kadau's Los Alamos Webpage

Molecular Dynamics → Continuum Mechanics

Long MD simulations represent nanoseconds .

Large MD simulations represent microns .

Continuum Mechanics uses **Finite Elements** rather than Atoms .



$$\dot{\rho} = -\rho \nabla \bullet \mathbf{v}$$

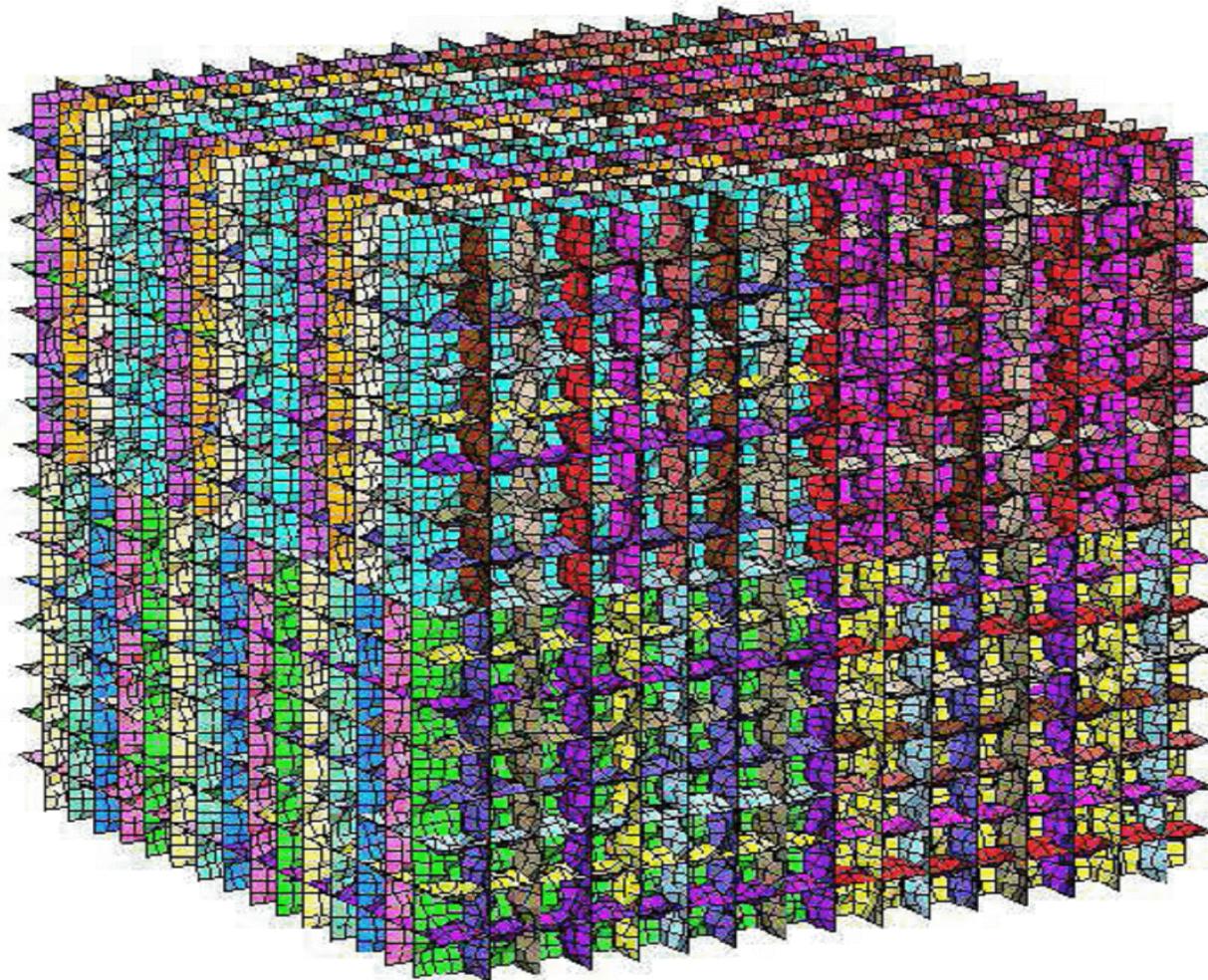
$$\rho \dot{\mathbf{v}} = -\nabla \bullet \mathbf{P} + \rho \mathbf{g}$$

$$\rho \dot{\mathbf{e}} = -\nabla \mathbf{v} : \mathbf{P} - \nabla \bullet \mathbf{Q}$$



The **Finite Elements** can be either
Eulerian (fixed) or **Lagrangian** (comoving) .

Auxetic Material Modeled with 208,896 Elastic-Plastic Shell Finite Elements .



Simpler Way: use **Particles governed by SPAM !
SPAM uses $w(r)$ to sum **Particle Contributions****

The “weight function” $w(r < h)$ smoothes out the **particles** .
This provides also very smooth averaged field variables .

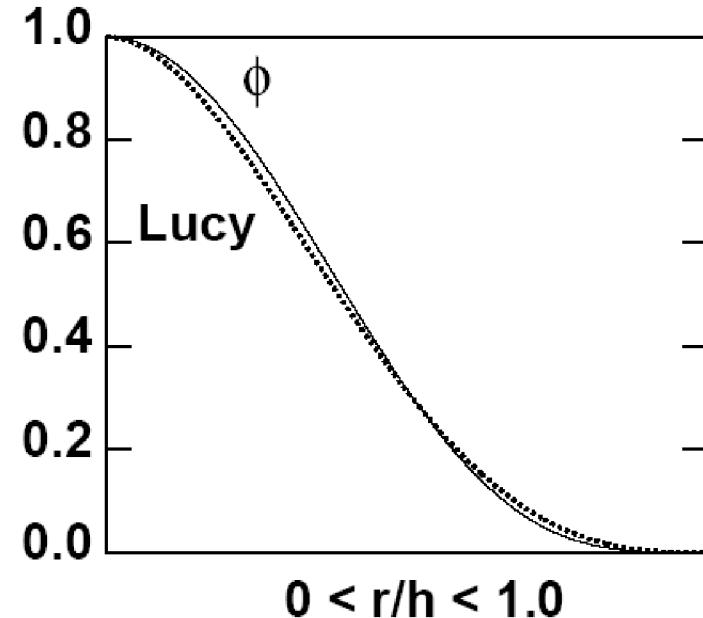
$$\rho_i = \sum_j m w_{ij}$$

$$\rho_i F_i = \sum_j m F_j w_{ij} \Rightarrow \nabla(\rho_i F_i) = \sum_j m F_j \nabla w_{ij}$$

$$\dot{v}_i = -\sum_j m [(\mathbf{P}/\rho^2)_i + (\mathbf{P}/\rho^2)_j] \bullet \nabla_i w_{ij} + g$$

Motion governed by *Ordinary* Differential Equations !

Lucy's weight function, $w = 1 - 6r^2 + 8r^3 - 3r^4$.



Short-range repulsive potential, $\phi = (1 - r^2)^4$.

SPAM can be isomorphic to Molecular Dynamics

Consider the case where P and ρ vary slowly in space .

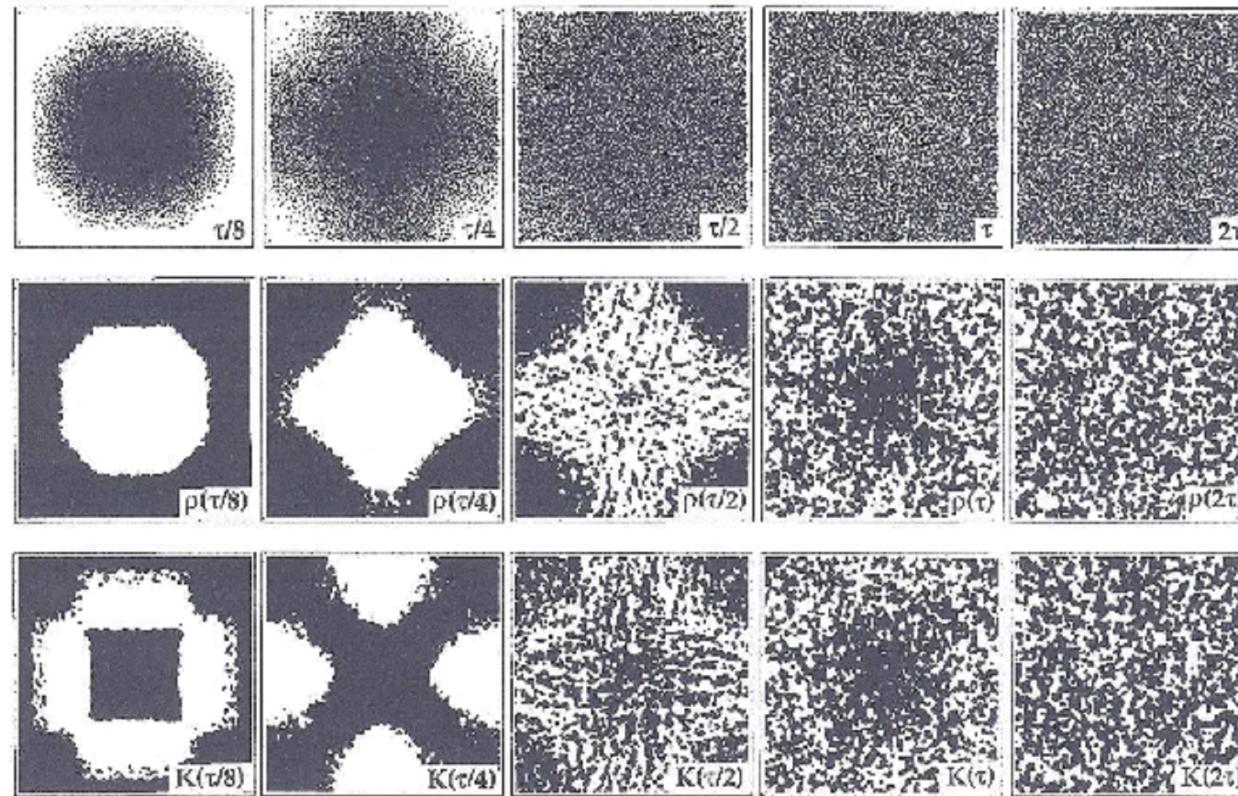
Consider the case of a gamma law fluid with $\gamma = 2$.

$$\dot{v}_i = - \sum_j m [(P/\rho^2)_i + (P/\rho^2)_j] \bullet \nabla_i w_{ij} \approx$$

$$\dot{v}_i \propto -\nabla_i w_{ij} \text{ [here SPAM is just like MD]}$$

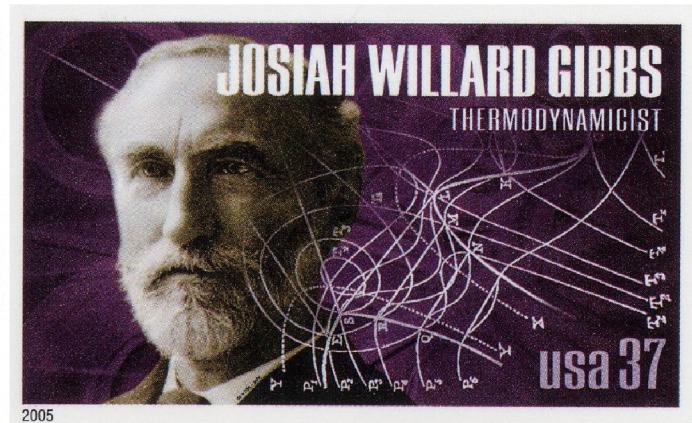
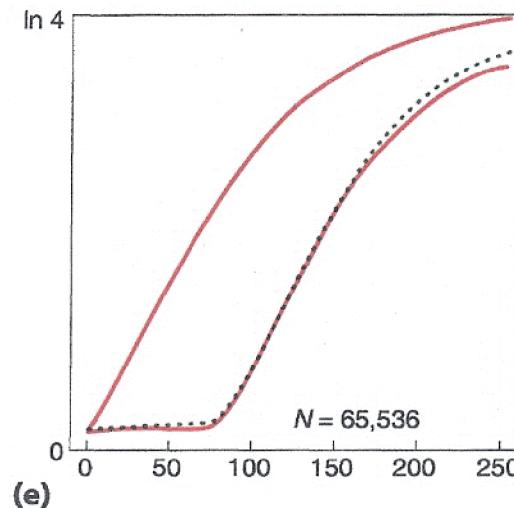
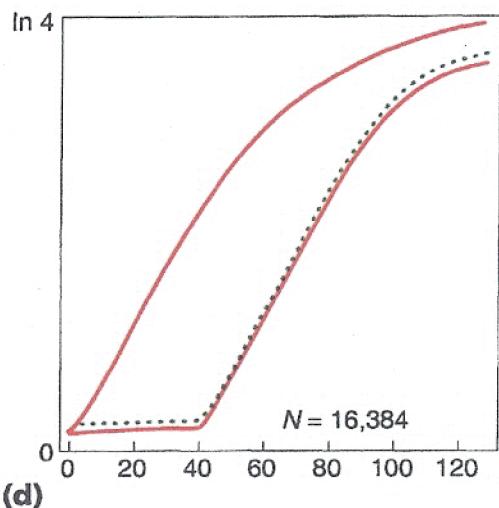
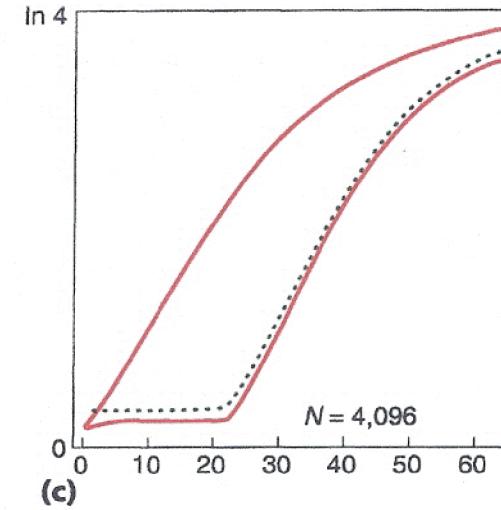
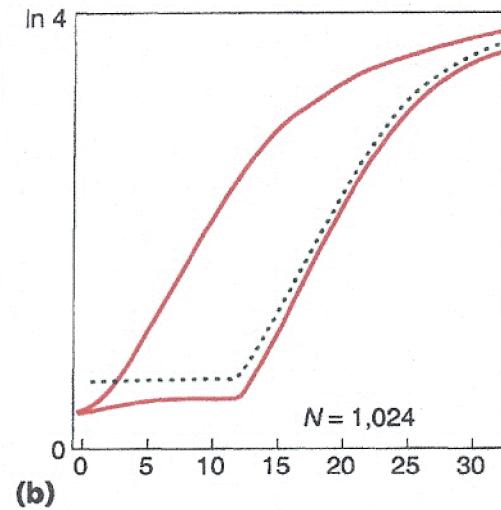
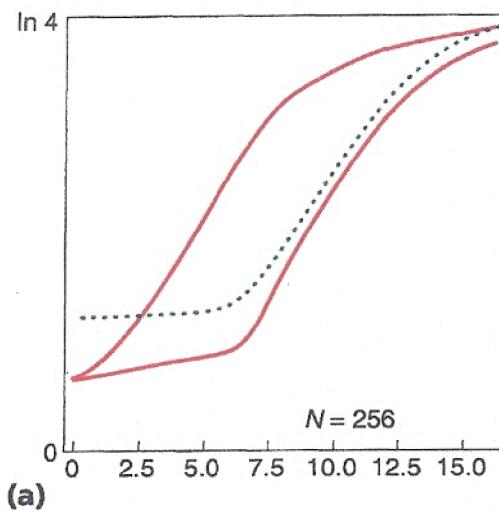
Consider a pedagogical example: Free Expansion .

Fourfold Free Expansion Simulation : **SPAM** calculation of ρ and K contours .

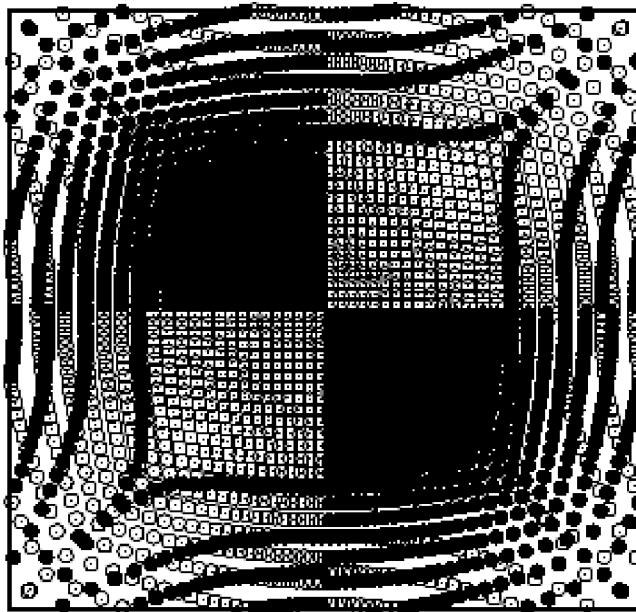
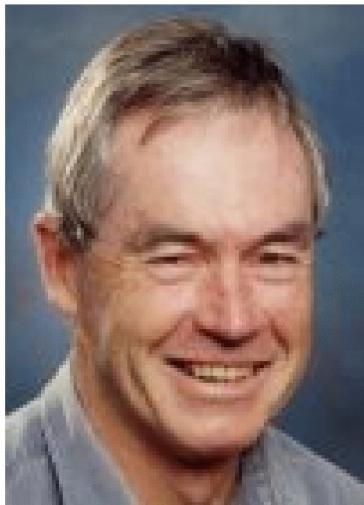


Equilibrates at the Sound Speed .

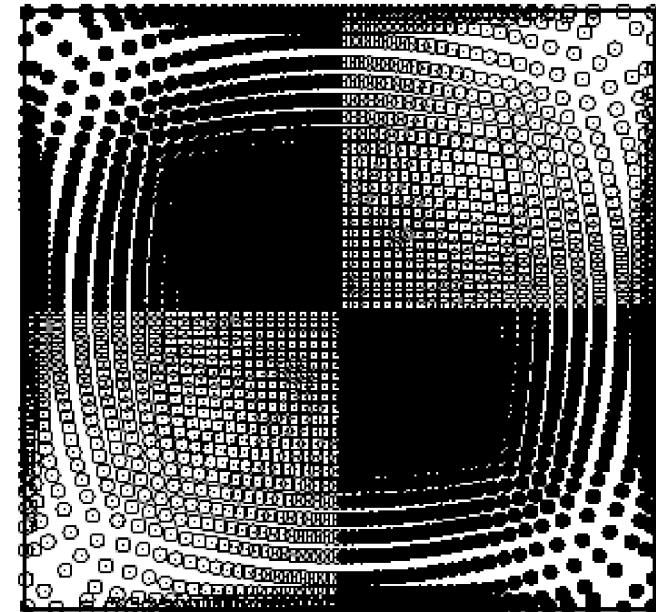
Fourfold Free Expansion Simulation : SPAM calculation of fluctuations in v .



For Continuum Simulations use Monaghan's Trick to Limit Fluid-Fluid Interpenetration .



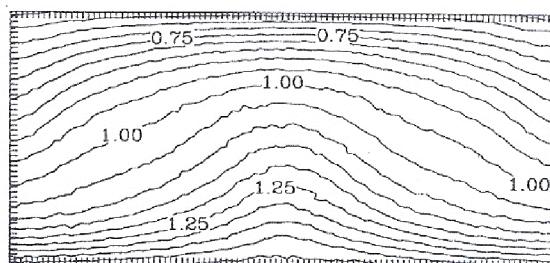
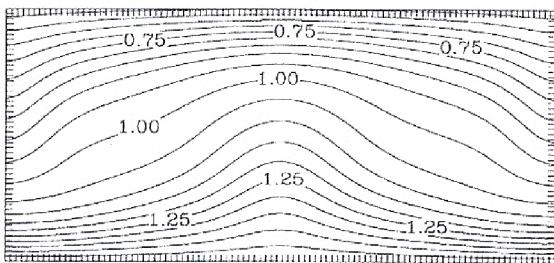
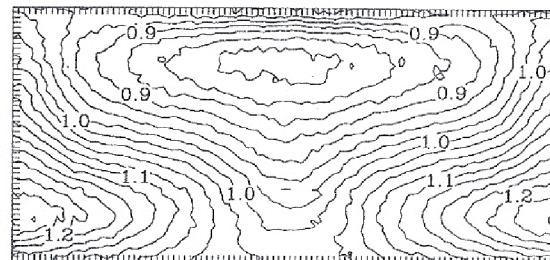
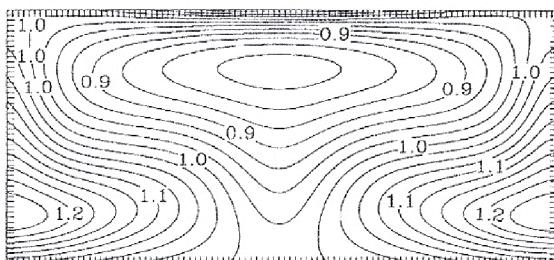
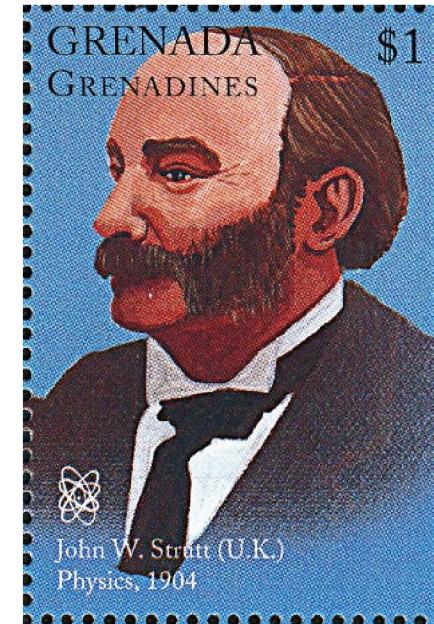
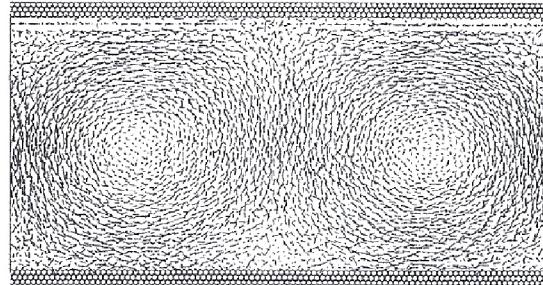
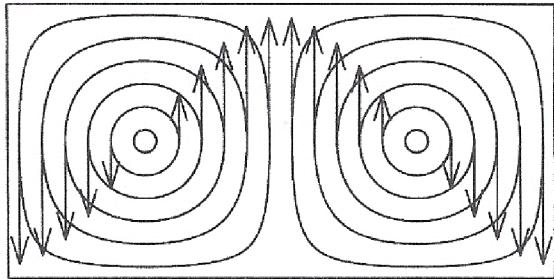
$$-25 < (x, y) < +25$$



$$-25 < (x, y) < +25$$

$$\{\dot{r} = v\} \Rightarrow \{\dot{r}_i = v_i + \sum_j m(v_j - v_i) w_{ij} / \rho_{ij}\}$$

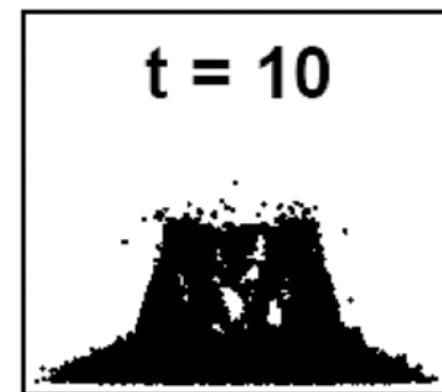
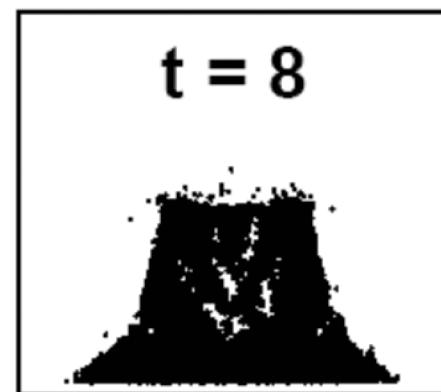
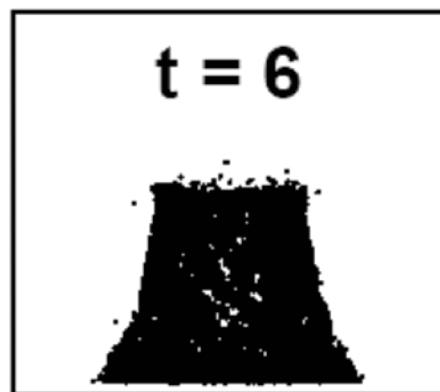
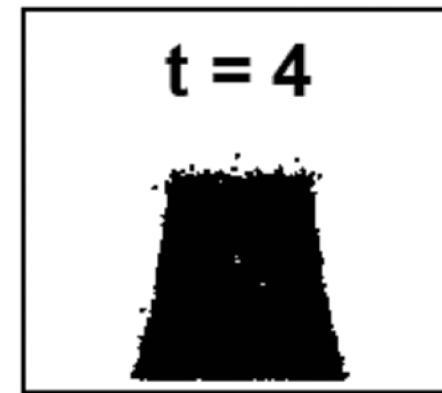
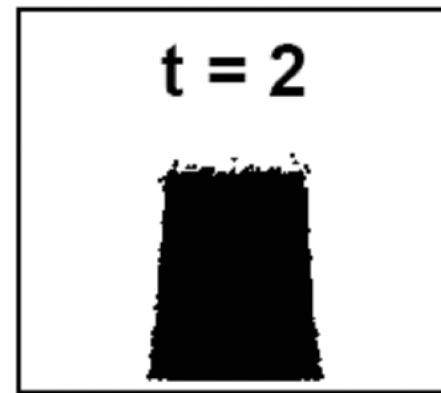
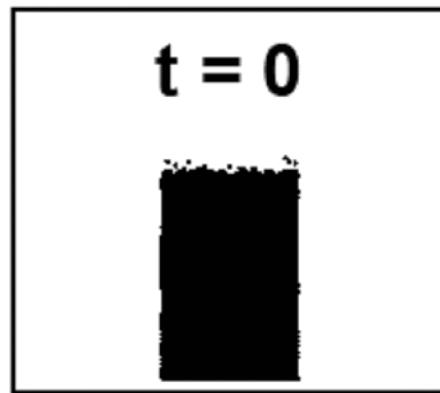
SPAM for Rayleigh-Benard Convective Instability



$\leftarrow N = 5000$
 $0.5 < T < 1.5$

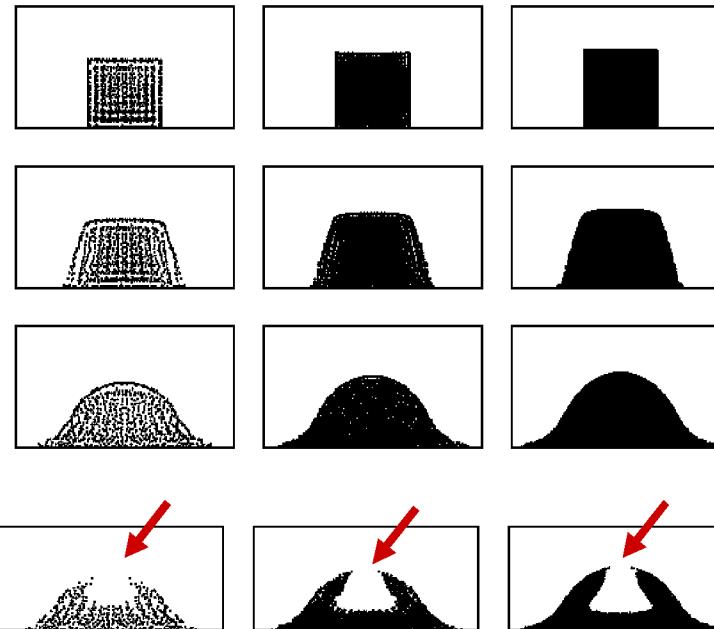
Continuum Mechanics versus SPAM, from Oyeon Kum's Thesis .

Column Collapse via Molecular Dynamics



$$\phi = (2 - r^2)^8 - 2(2 - r^2)^4 \text{ for } r^2 < 2.$$

Column Collapse via SPAM with $P = \rho^3 - \rho^2$



$N = 640$;
 $N = 2560$;
 $N = 10240$.

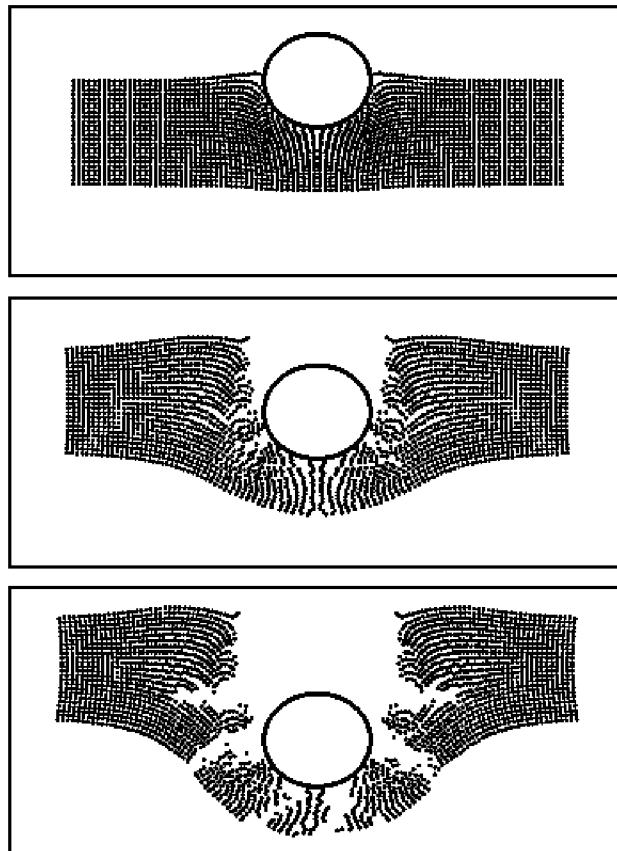
Tensile Regions

Surface potential and Core potential are both included:

$\Phi \propto (\nabla \rho)^2$ for minimizing surface .

$\Phi \propto \Sigma(\sigma^2 - r^2)^4$ for avoiding string formation .

Penetration using SPAM With Tensile Failure Model



$$\rho = \sqrt{4/3} ; v = 4 ;$$
$$Y = \eta/100 ; \sigma = \eta/2 ;$$
$$\eta = \lambda = \sqrt{24}.$$

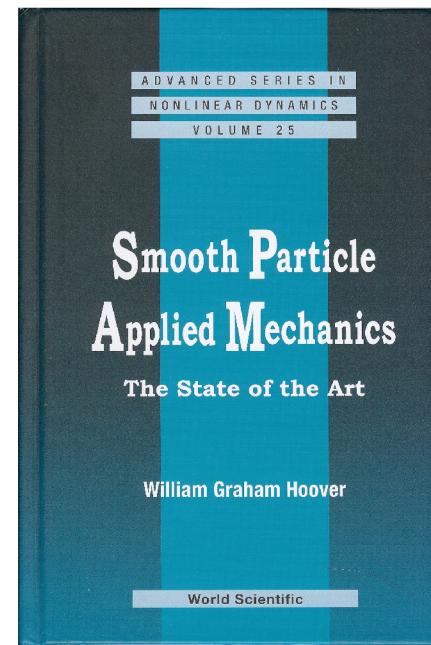
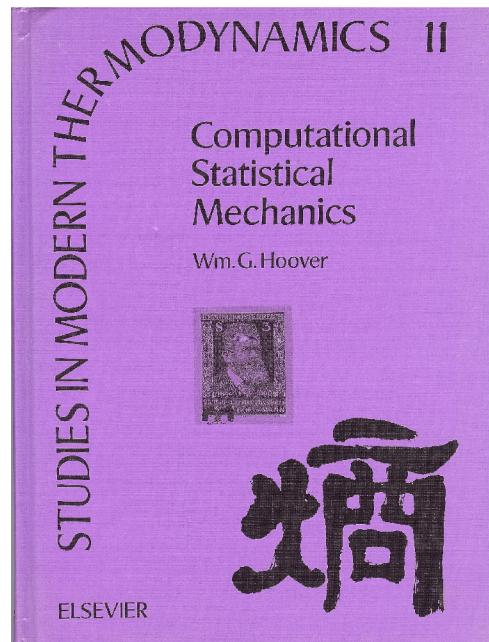
$-60 < x < +60 ; t = 4, 12, 20$

“What to do?” [Lenin]

Encourage Students to Implement Particle Methods .

Molecular Dynamics is Classical, while Real Life is not .

SPAM has a need for documented failure models .



<http://williamhoover.info>